

Singularities driven turbulence

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Motivation-context

- Role (and in fact existence) of singularities (or at least extreme events) in turbulence
- Singularities in Navier-Stokes equations (back to Leray 1934) are still a question of debates
- Extreme events (vorticity collapses) have been predicted and observed numerically (Siggia-Pumir 1987, Brachet et al 1992 for instance) and experimentally (Meneveau-Sreenivasan 1987). They are expected to be responsible of the intermittence in turbulence
- Difficult to handle in fluid flows: can we investigate a simpler model where we know the singularities (hard turbulence 89-90, burgulence Bec-Frisch 2000')?

Intermittent fluctuations of the energy dissipation (Meneveau & Sreenivasan, Nuclear Physics B, 1987)

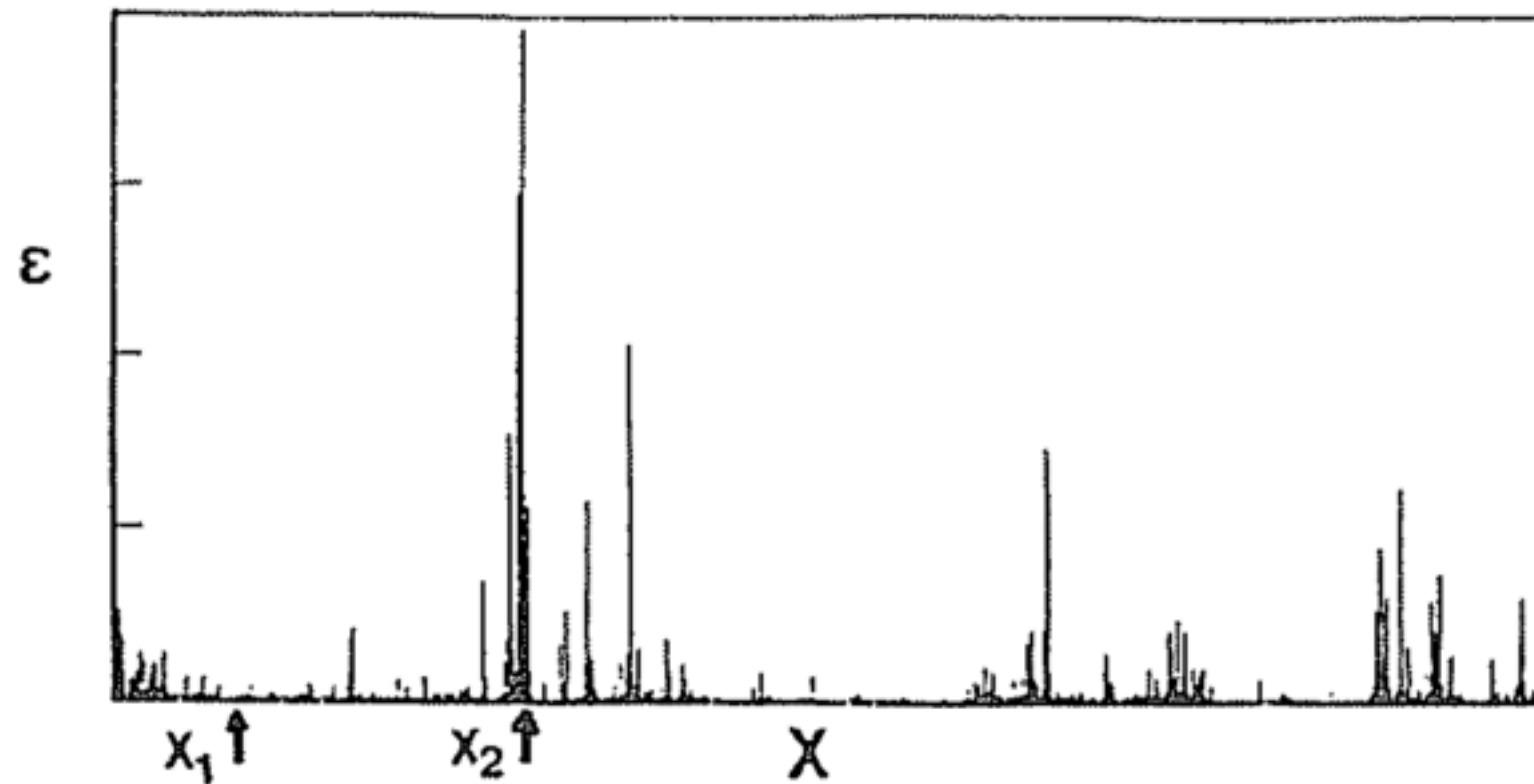


Figure 1a. Typical time trace of $(\partial u_1 / \partial t)^2$, representative of the rate of dissipation of turbulent kinetic energy.

Focusing NLS

$$i\frac{\partial\psi}{\partial t} = -\frac{\alpha}{2}\nabla^2\psi - g|\psi|^{2n}\psi$$

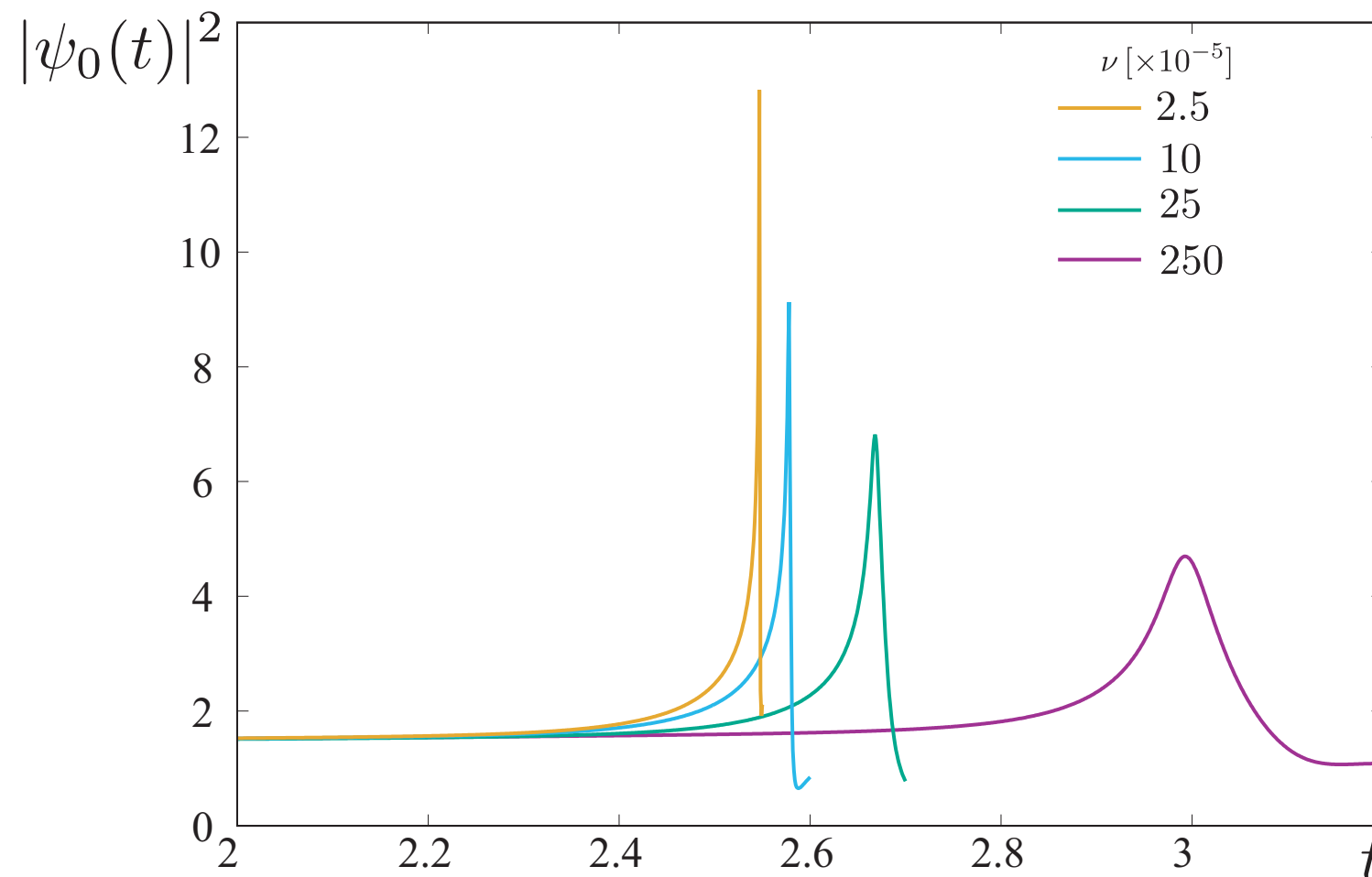
- $\psi(\mathbf{x}, t)$ is a complex field
- the NLS equation is a model for BEC, shallow-water and nonlinear optics.
- Hamiltonian structure and mass conservation

$$N = \int |\psi|^2 d\mathbf{x}$$

$$H = \int \left(\frac{\alpha}{2} |\nabla\psi|^2 - \frac{g}{n+1} |\psi|^{2(n+1)} \right) d\mathbf{x}$$

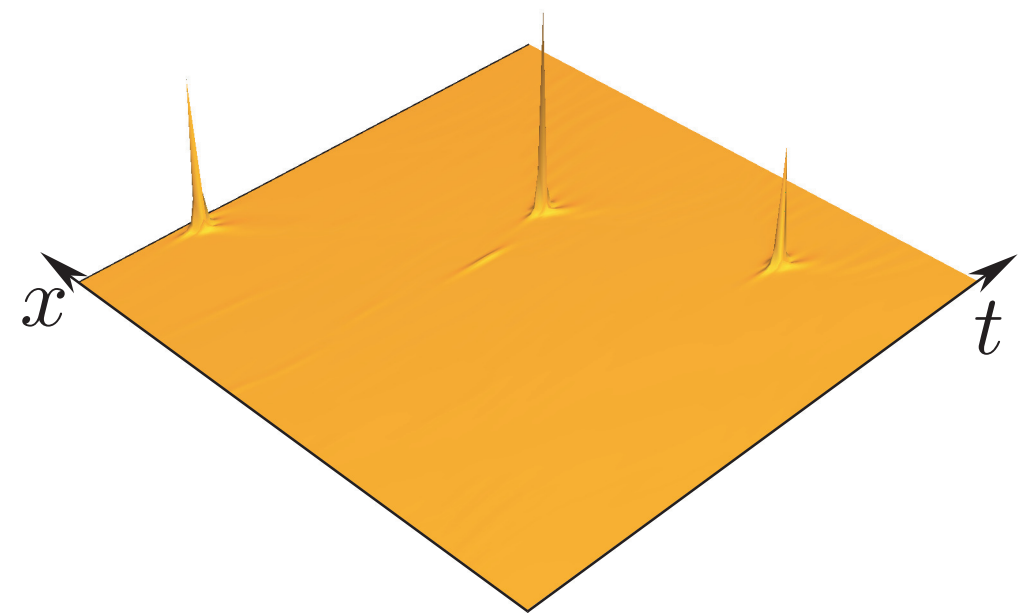
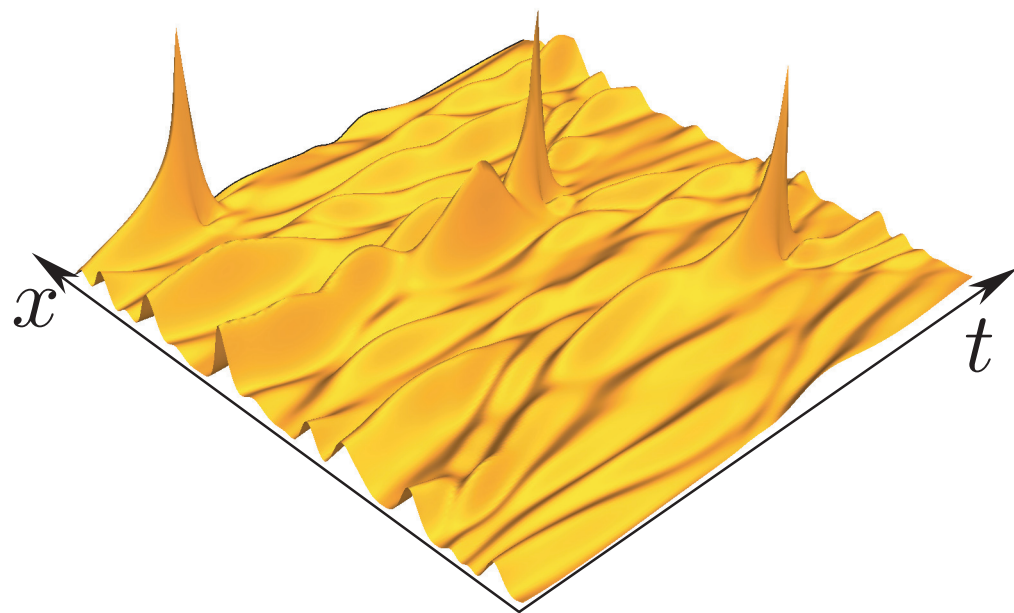
Singularity in NLS

- for the focusing NLS ($g=1$), the dynamics exhibits finite time and space singularities for $nd>2$ (see for instance Le Mesurier et al 1988)
- The singularity is due to the non positive defined Hamiltonian
- this finite time singularity is suppressed in the presence of dissipation



$$i\frac{\partial\psi}{\partial t} = -\frac{\alpha}{2}\nabla^2\psi - g|\psi|^{2n}\psi - i\nu\Delta^2\psi + f_{k_0}(\mathbf{x}, t).$$

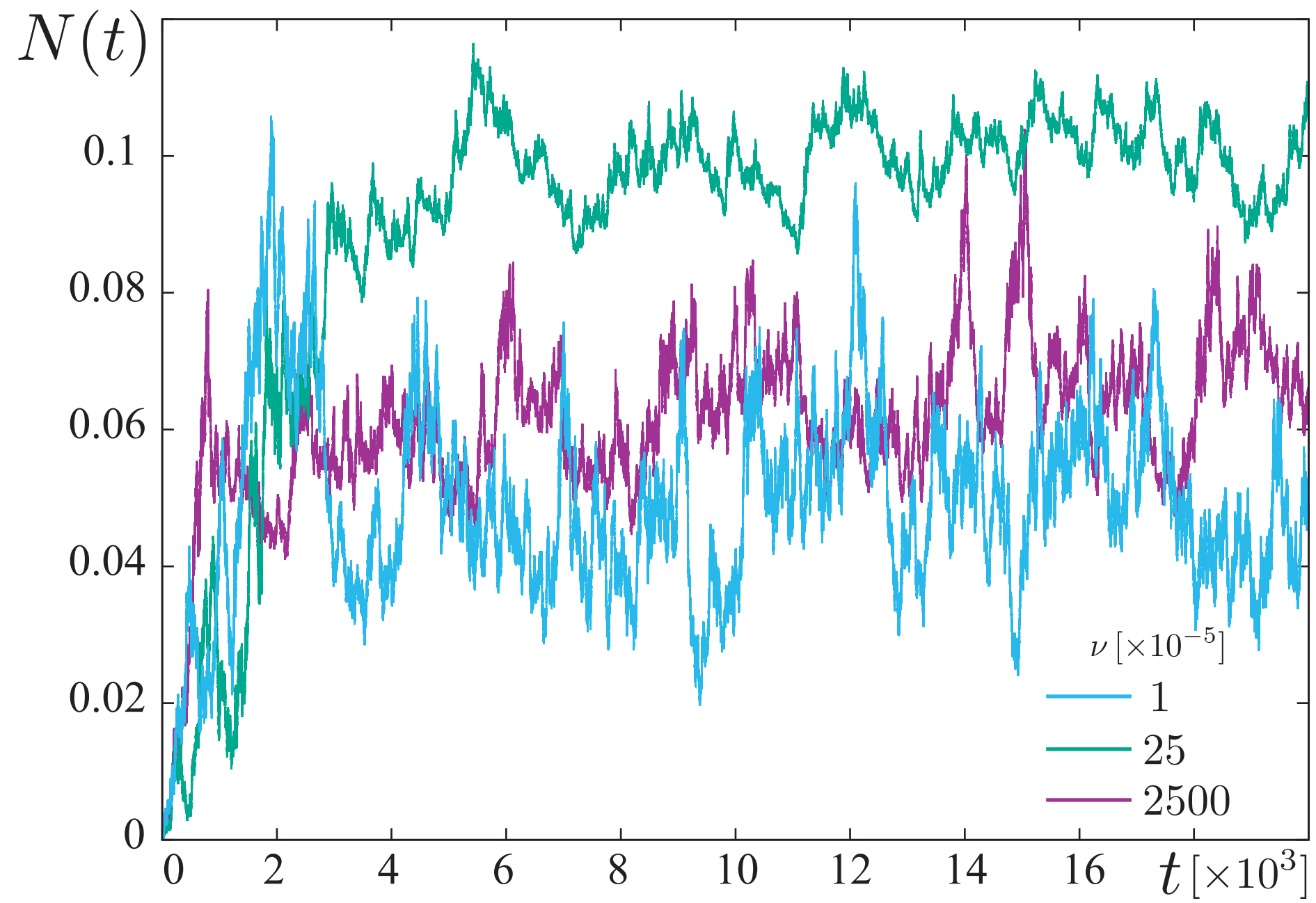
- with injection at large scale: turbulence of singularity or collapses (Dyachenko et al 1992, M. Bartuccelli et al 1989 & 1990)
- important difference: here the mass (positive definite) is the pertinent quantity for investigating turbulence
- wave turbulence (often observed in NLS equations) would suggest *inverse* cascade of mass and direct of energy
- focus here on the 1D case with $n=3$ (work in progress in 2 and 3D)
- numerical simulations

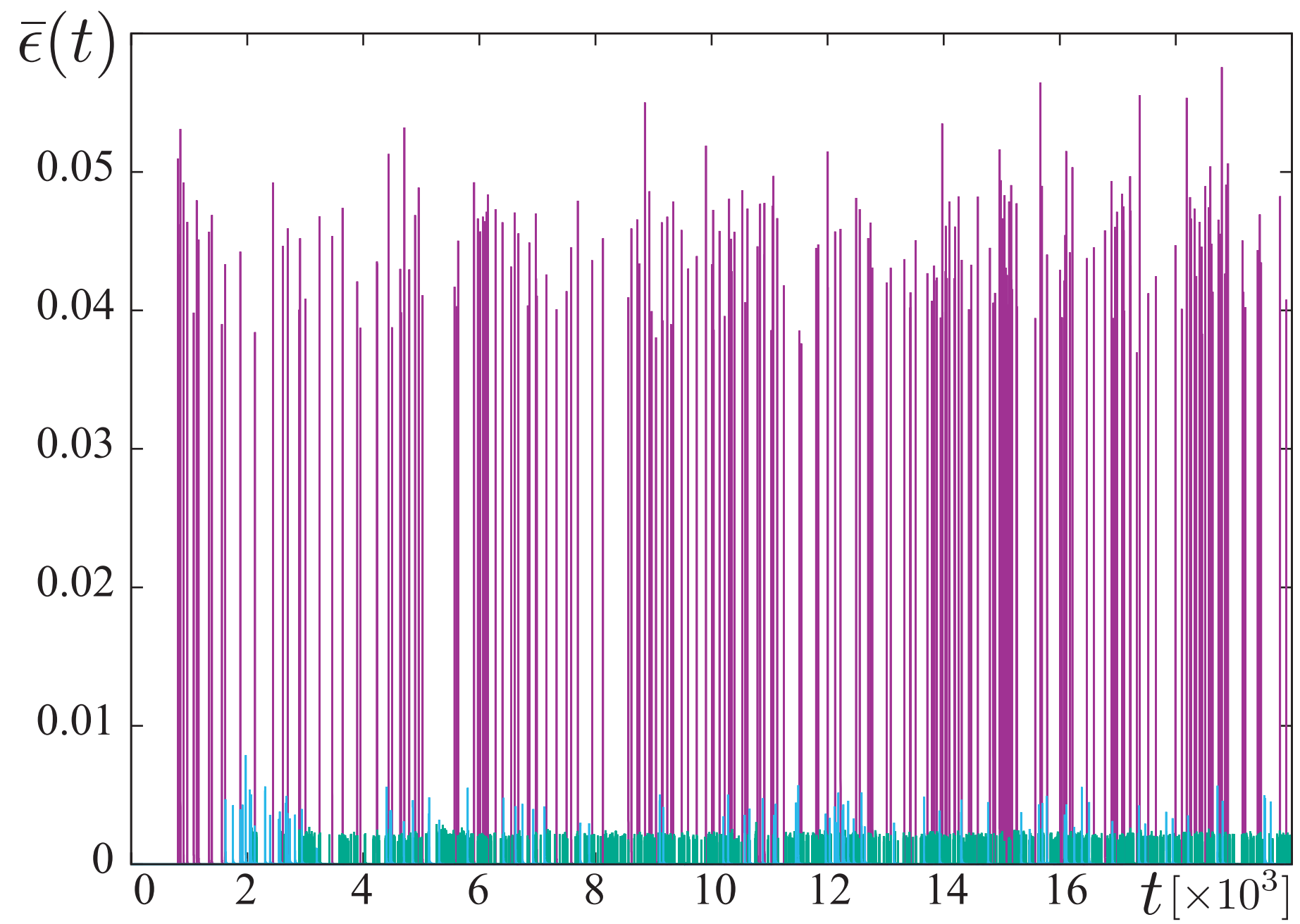


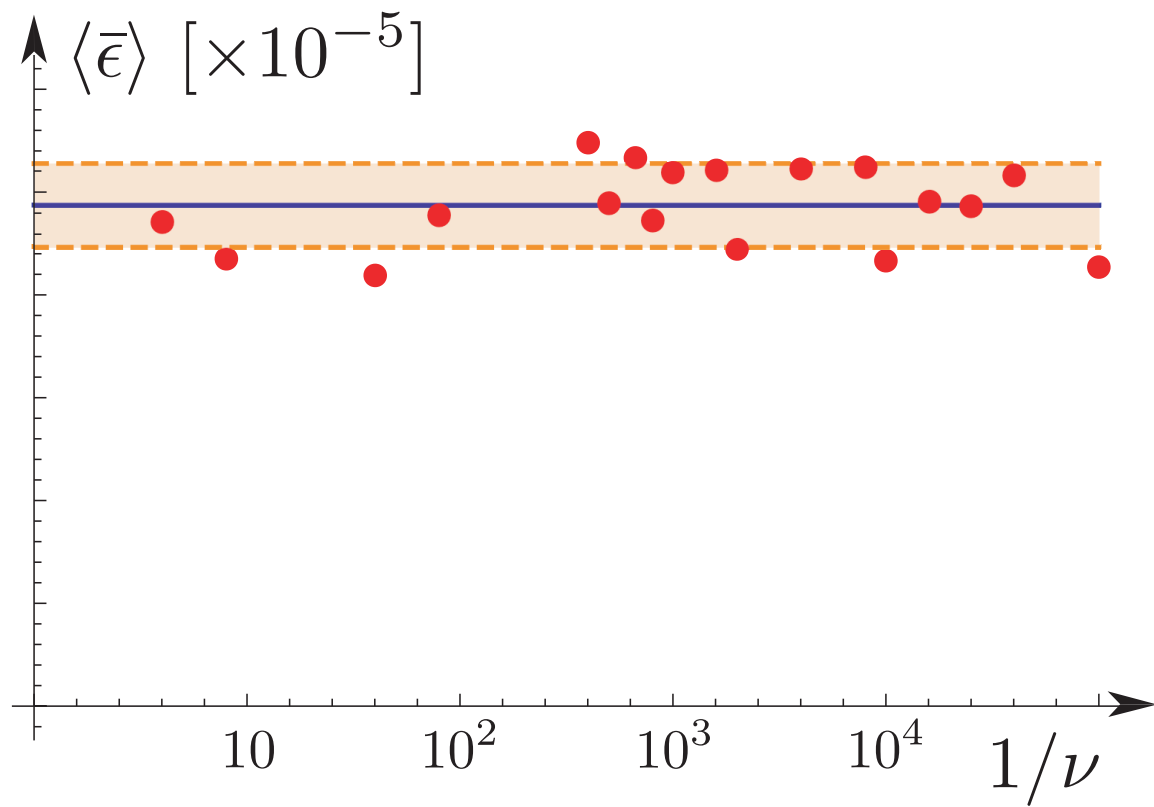
- « singularity » collapse or peak is followed by wave emission
- « dissipation » of mass is concentrated at short scale on the peaks:

$$\frac{dN}{dt} = -2\nu \int |\Delta\psi|^2 d^D\mathbf{x} + i \int (\psi \bar{f}_{k_0} - \bar{\psi} f_{k_0}) d^D\mathbf{x}.$$

Varying only the viscosity



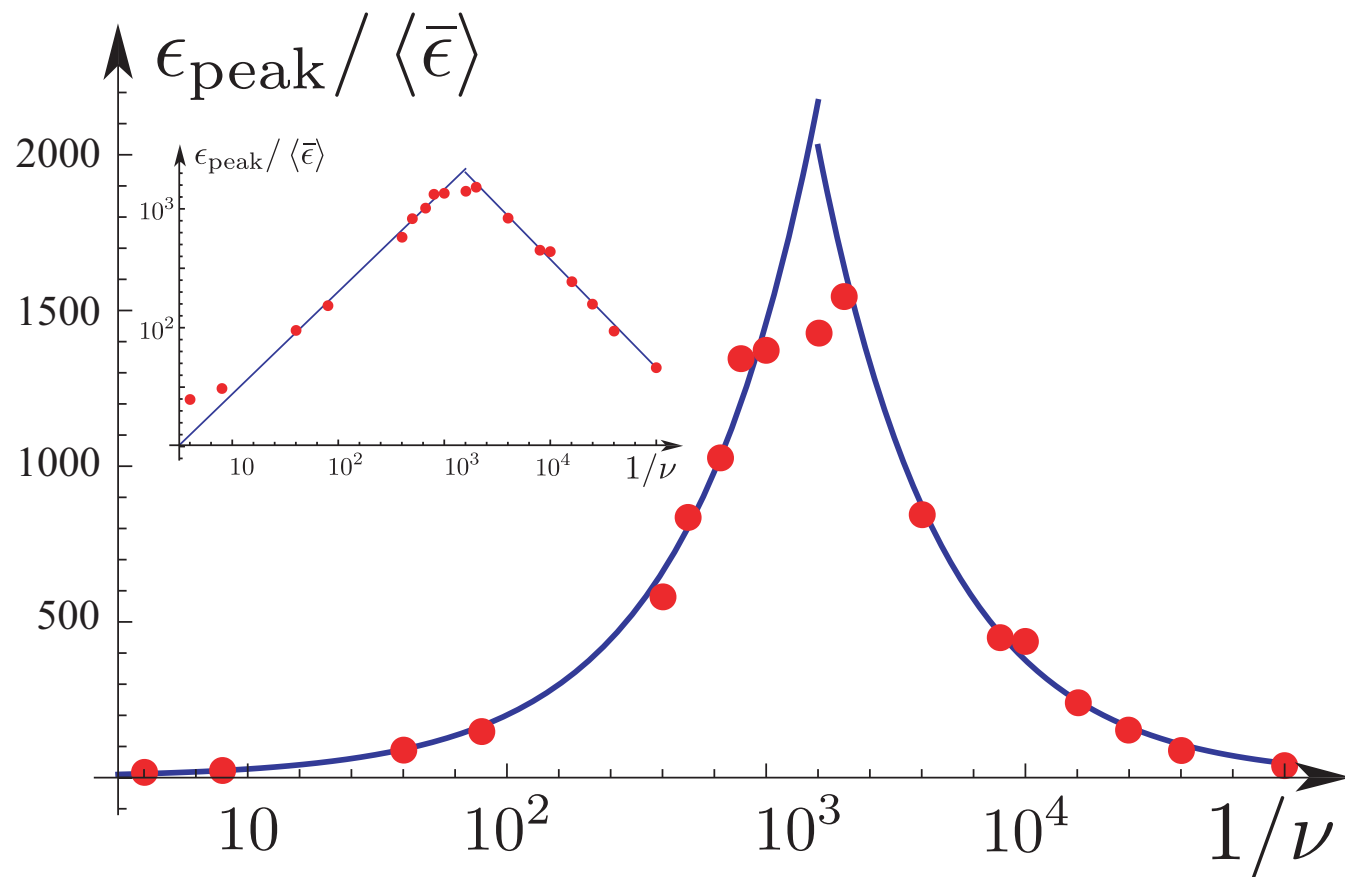




Anomalous
dissipation?

Warning:

$$\nu \times |\Delta\psi|^2$$

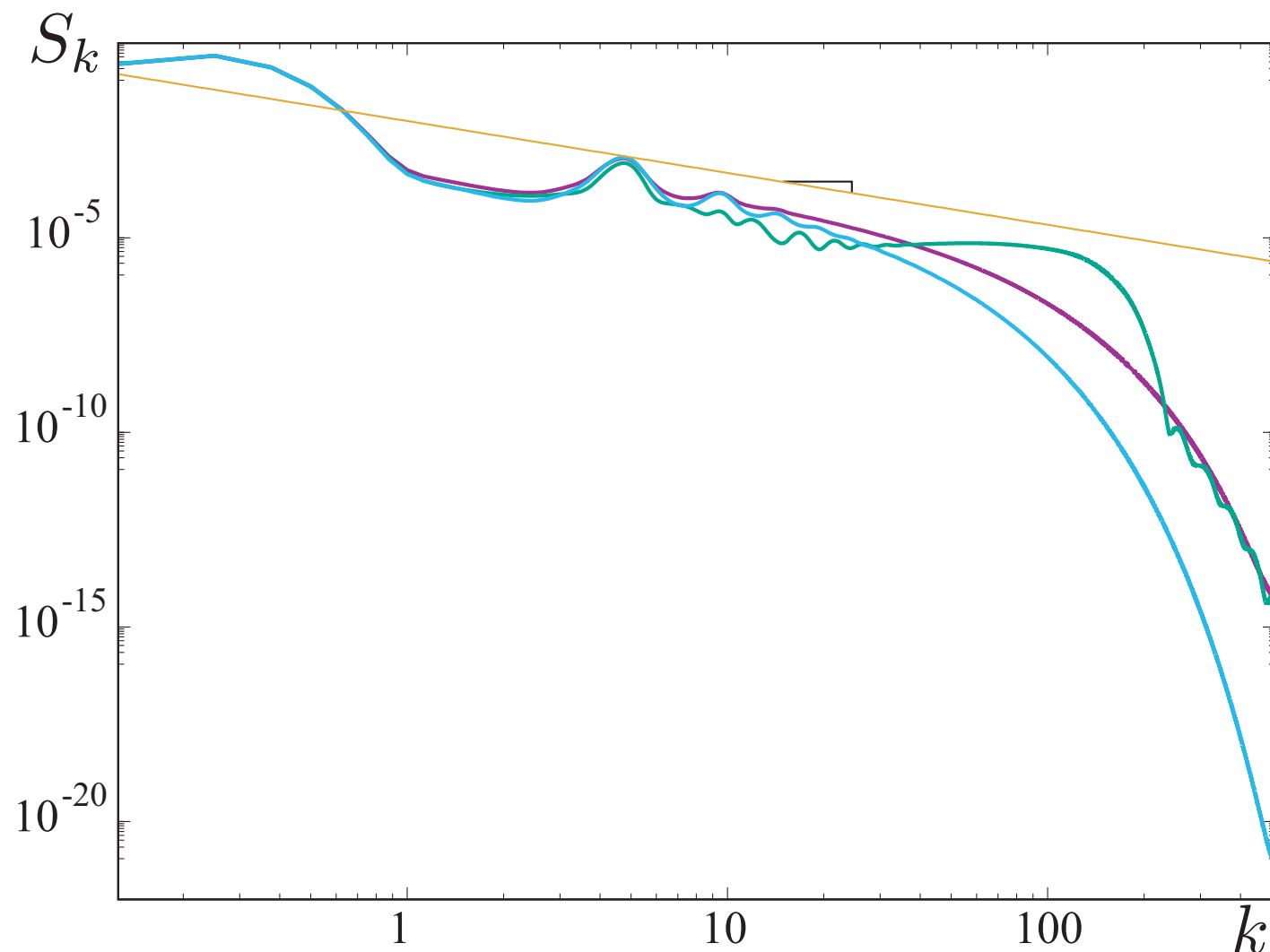


Spectrum

$$S_k(t) \equiv |\hat{\psi}_k|^2 + |\hat{\psi}_{-k}|^2$$

$$\frac{1}{L} \int |\psi|^2 d\mathbf{x} = \int |\hat{\psi}_k|^2 d\mathbf{k}$$

- Spectrum fluctuates at collapse



$$\frac{1}{t} \sim \frac{1}{x^2} \sim |\psi|^6$$

$$|\psi|^2 \sim x^{-2/3}$$

$$|\psi_k|^2 \sim k^{-4/3}$$

- spectrum of the self similar collapse

$$S_k \propto k^{-4/3}$$

- Transport equation for the spectrum

$$\frac{\partial S_k}{\partial t} = -\frac{\partial Q_k}{\partial k} - 2\nu k^4 S_k + F_k$$

- Need to consider averaged in time spectra for which we have

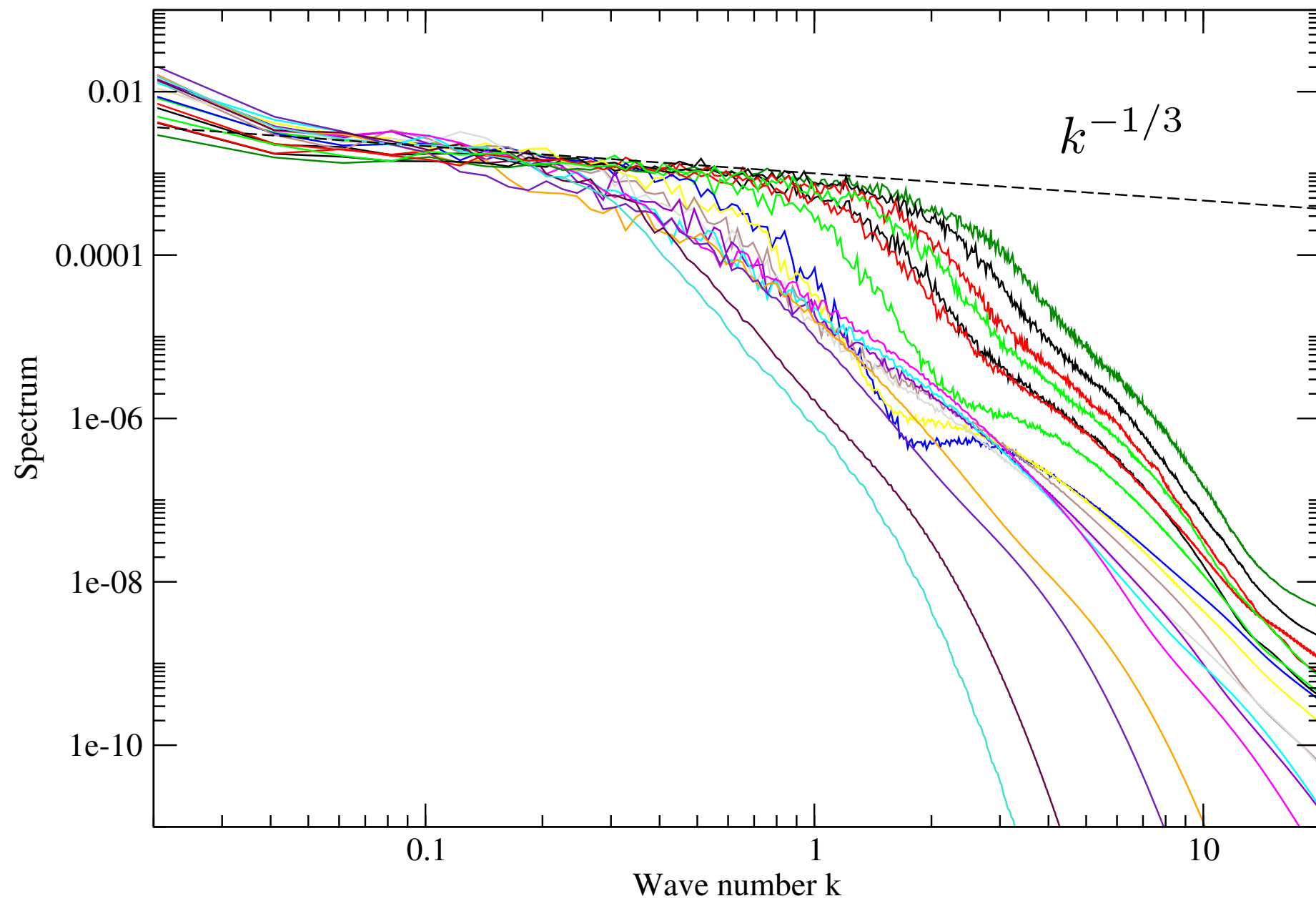
$$\frac{\langle \partial S_k \rangle}{\partial t} = 0$$

- In the inertial range we obtain also

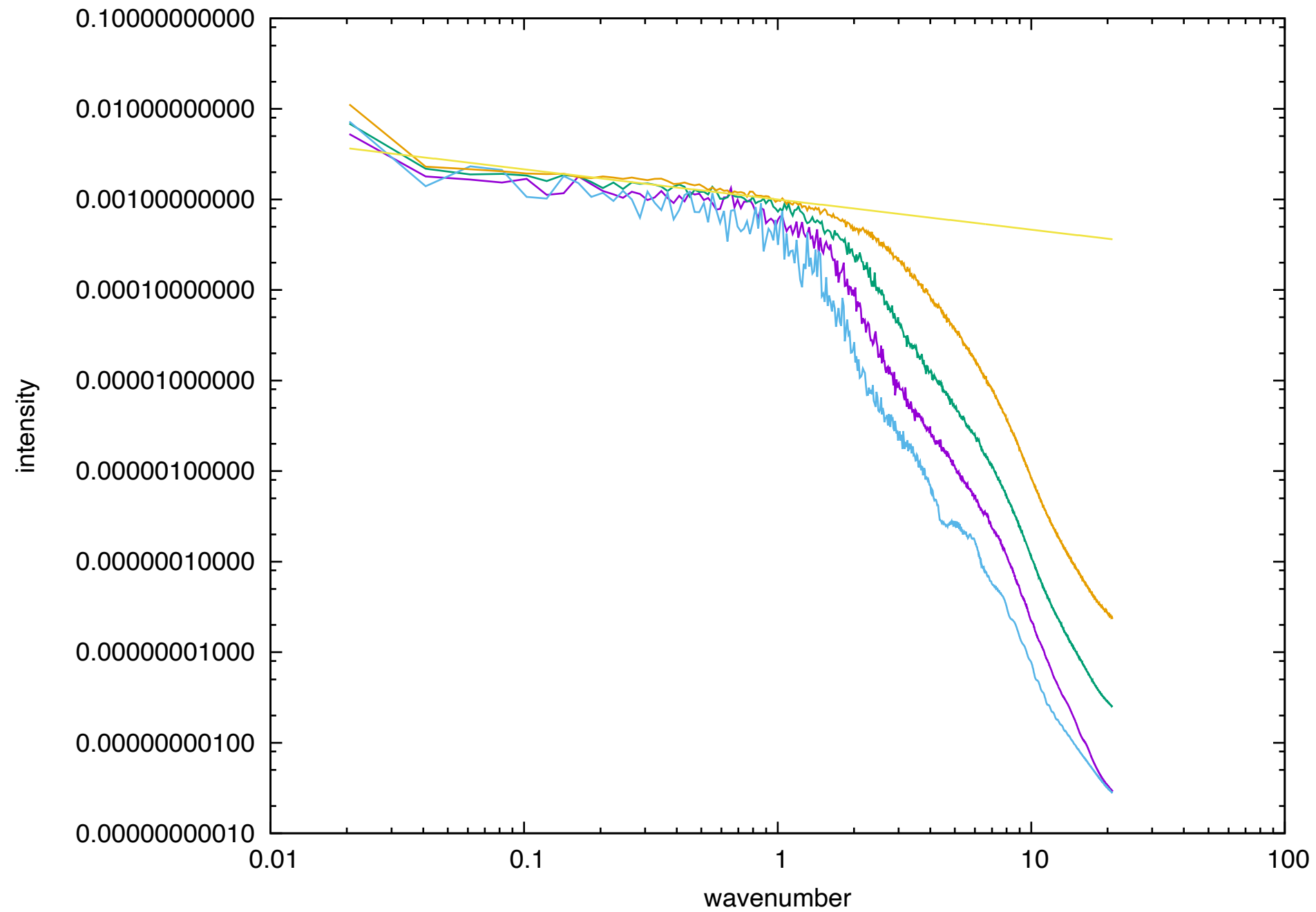
$$\langle Q_k \rangle = -2\nu \int_k^\infty k^4 \langle S_k \rangle dk \equiv \langle \epsilon \rangle$$

•

Spectrum expands in k as the viscosity decreases. Its amplitude seems independent of the injection rate



Varying the injection



Phillips spectrum?

Kolmogorov-like scaling analysis

$$[S_k] = \rho \ell \quad [\epsilon] = \rho \tau^{-1} \quad [\alpha] = \ell^2 \tau^{-1} \quad [g] = \rho^{-3} \tau^{-1}$$

$$\langle S_k \rangle = \frac{\langle \bar{\epsilon} \rangle}{\alpha k^3} F \left(\frac{\alpha k^2}{(g \langle \epsilon \rangle^3)^{\frac{1}{4}}} \right)$$

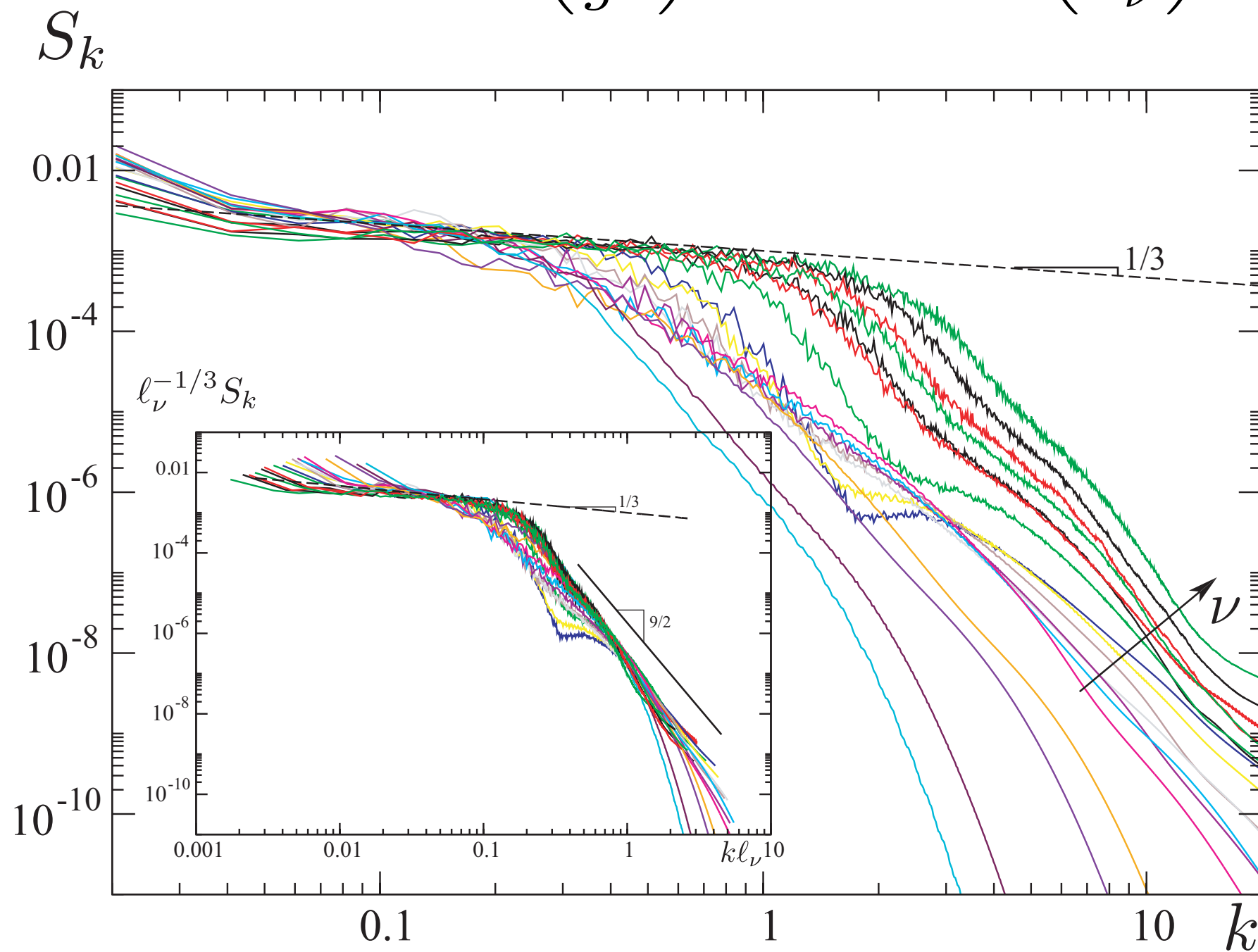
If we look for a solution independent of the injection rate

$$\langle S_k \rangle \propto \left(\frac{\alpha}{g^4} \right)^{1/3} k^{-1/3}$$

$$\text{Kolmogorov scale } \ell_\nu \sim \left(\frac{\alpha \nu^3}{g \bar{\epsilon}^3} \right)^{1/14} \quad k_\nu \sim \left(\frac{g \bar{\epsilon}^3}{\alpha \nu^3} \right)^{1/14}$$

Suggest the following self-similar scaling for the spectrum

$$\langle S_k \rangle = \left(\frac{\alpha}{g^4} \right)^{1/3} k_\nu^{-1/3} G \left(\frac{k}{k_\nu} \right)$$



Intermittency-structure functions

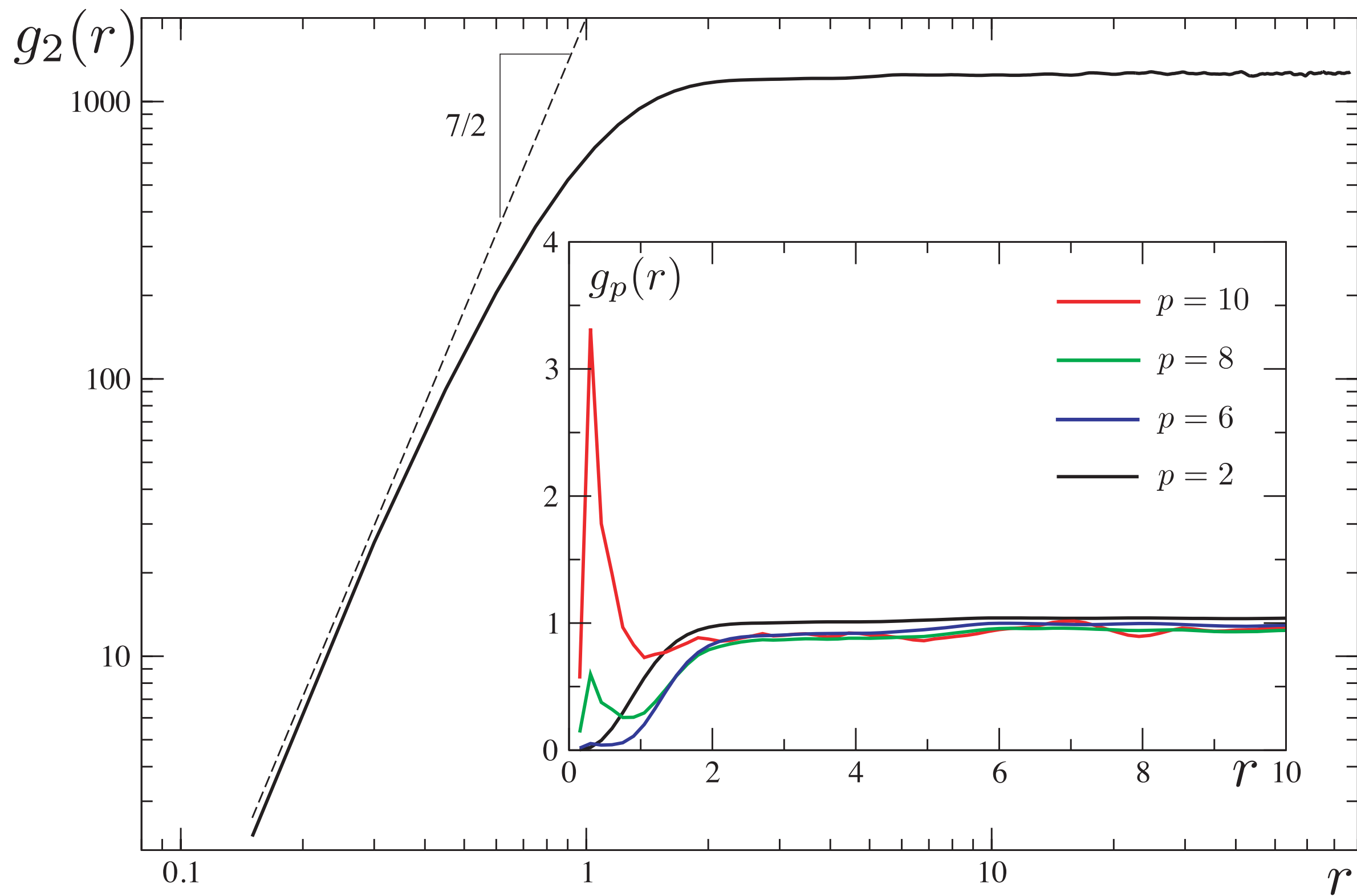
$$g_p(r) = \overline{|\psi(x+r) + \psi(x-r) - 2\psi(x)|^p}$$

p=2 can be deduced from the spectrum scalings

$$g_2(r) \sim r^{7/2} \quad \text{at short scales}$$

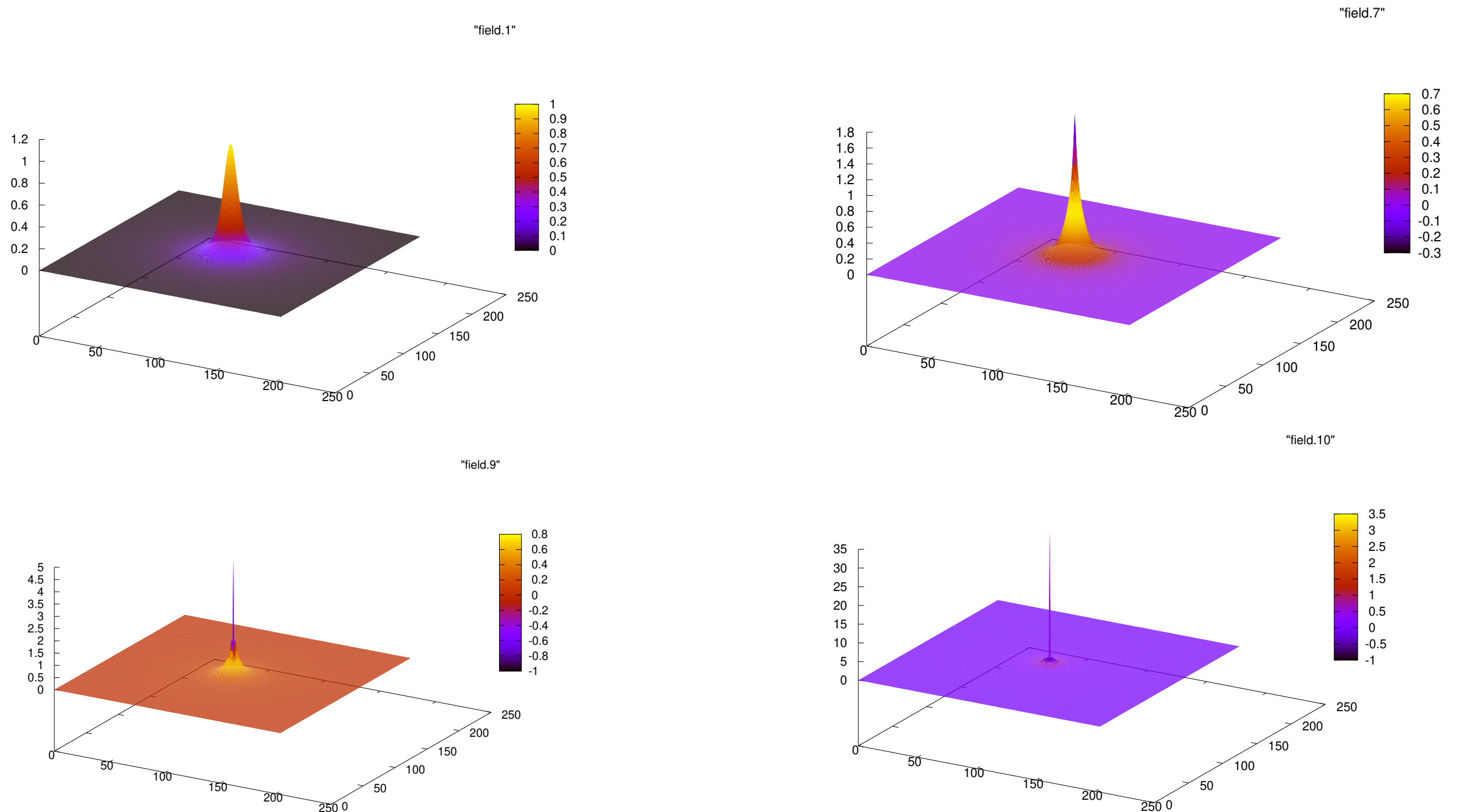
$$g_2(r) \sim r^{-2/3} \quad \text{inertial range}$$

High p's should witness the singularity at small scales

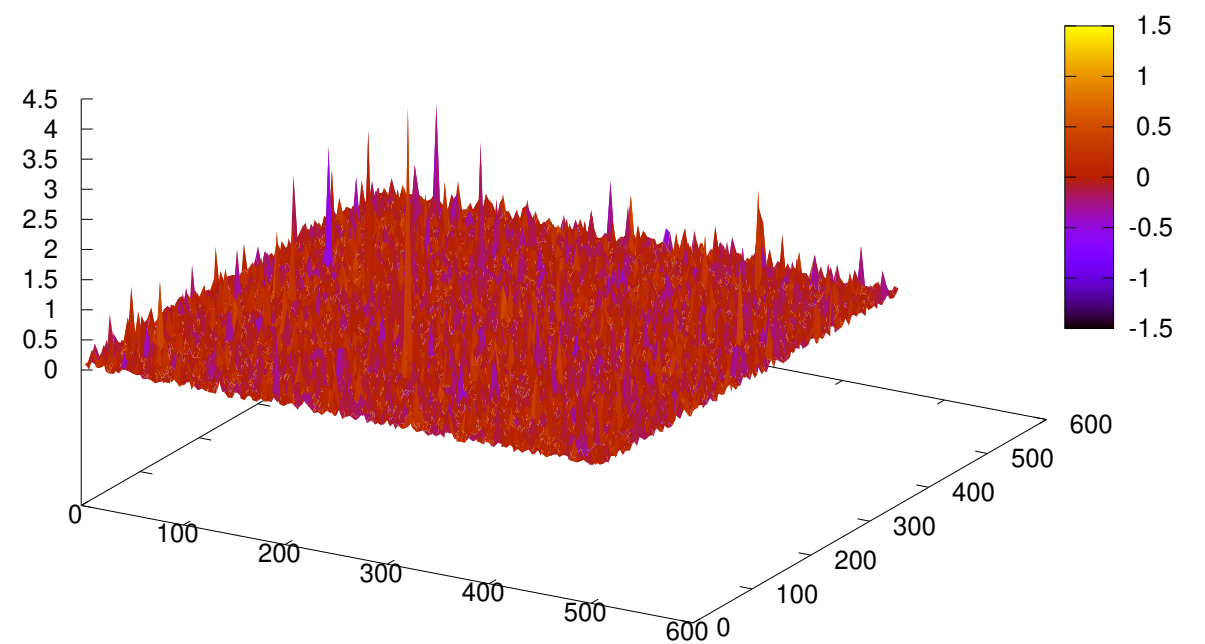
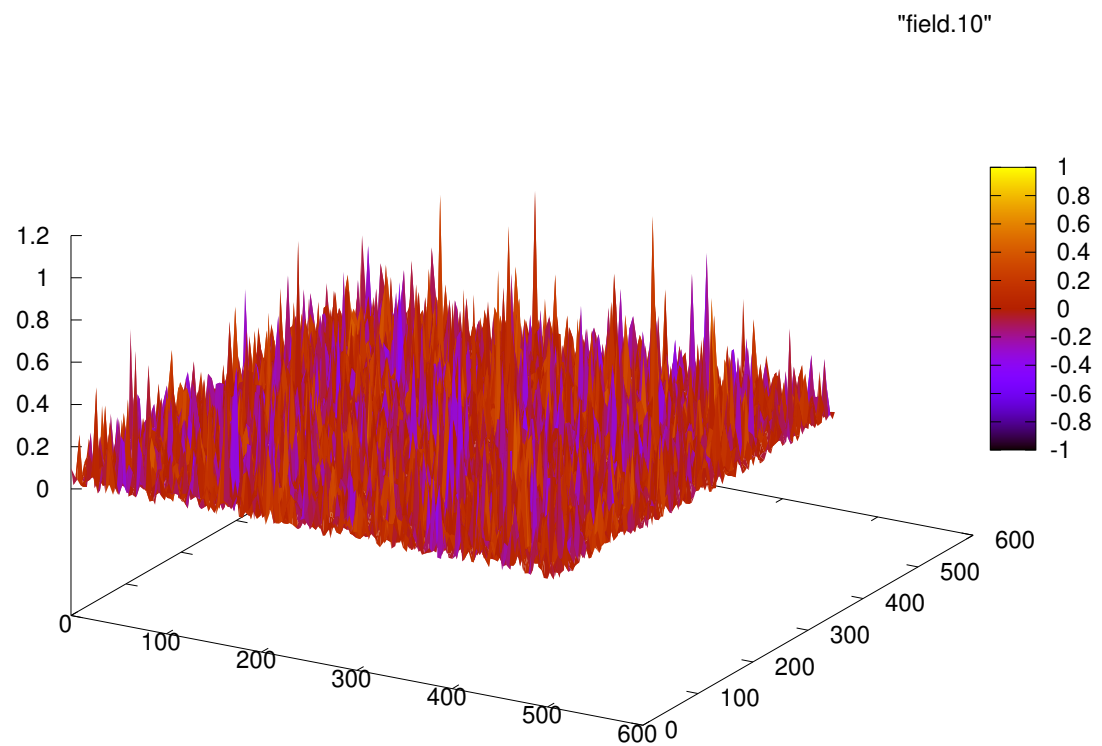
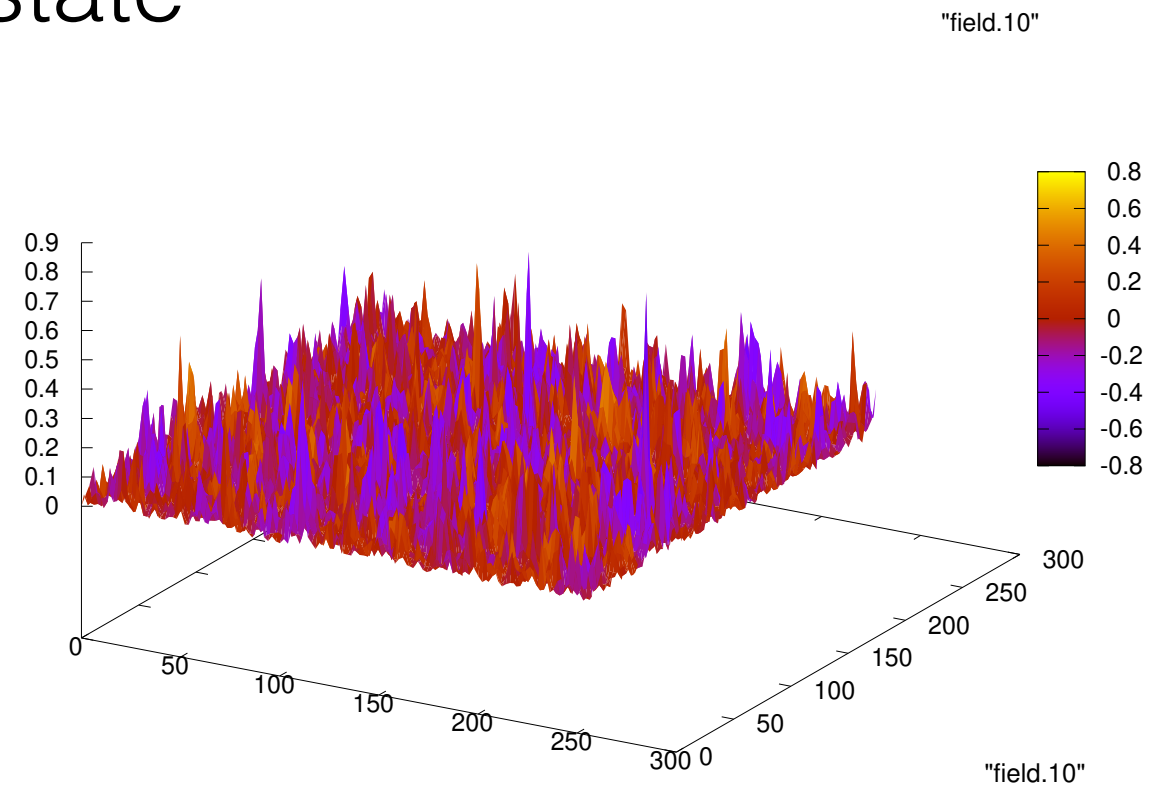
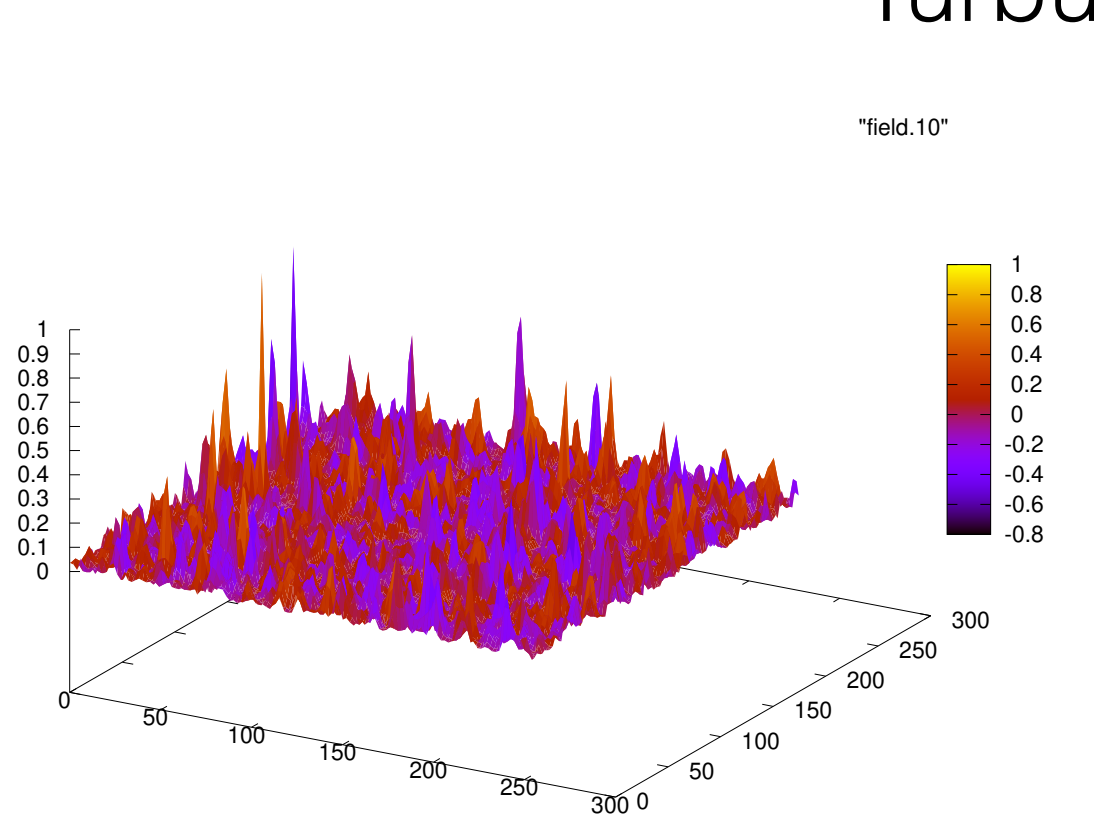


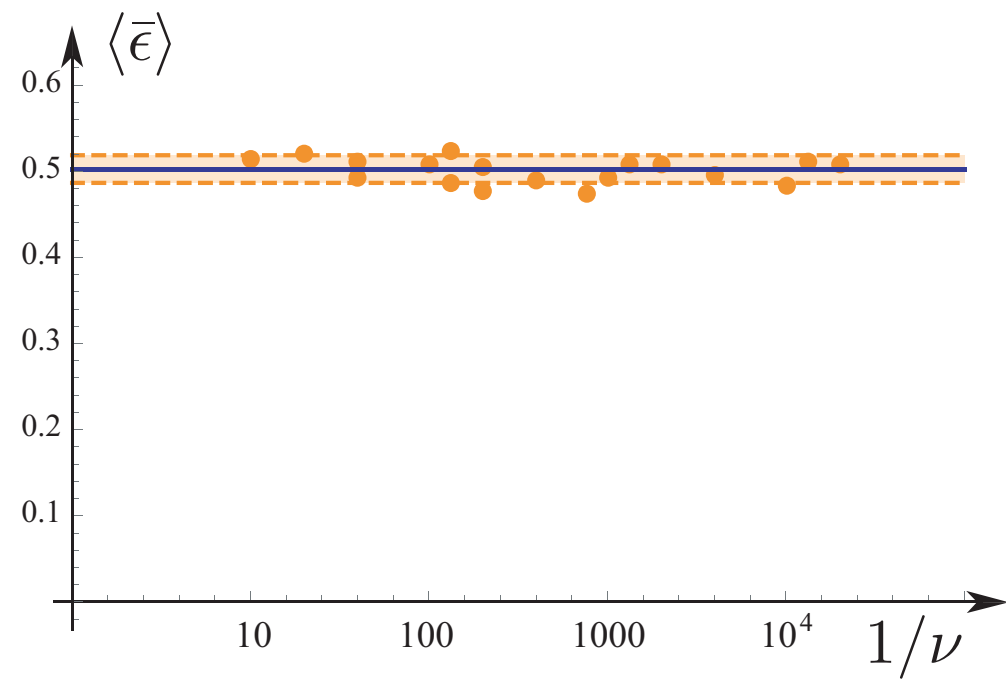
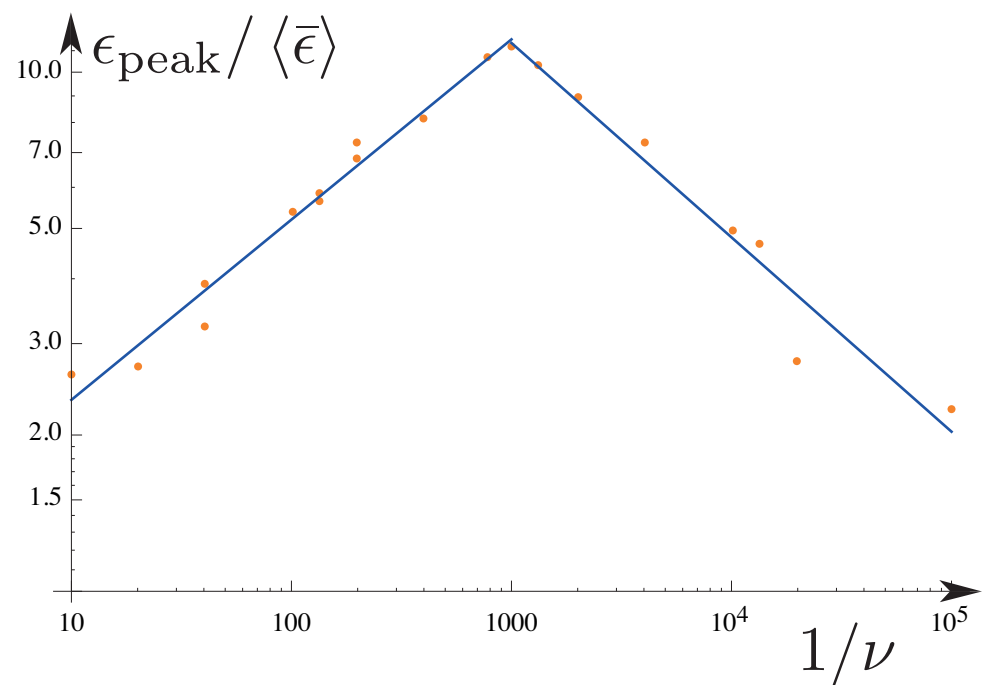
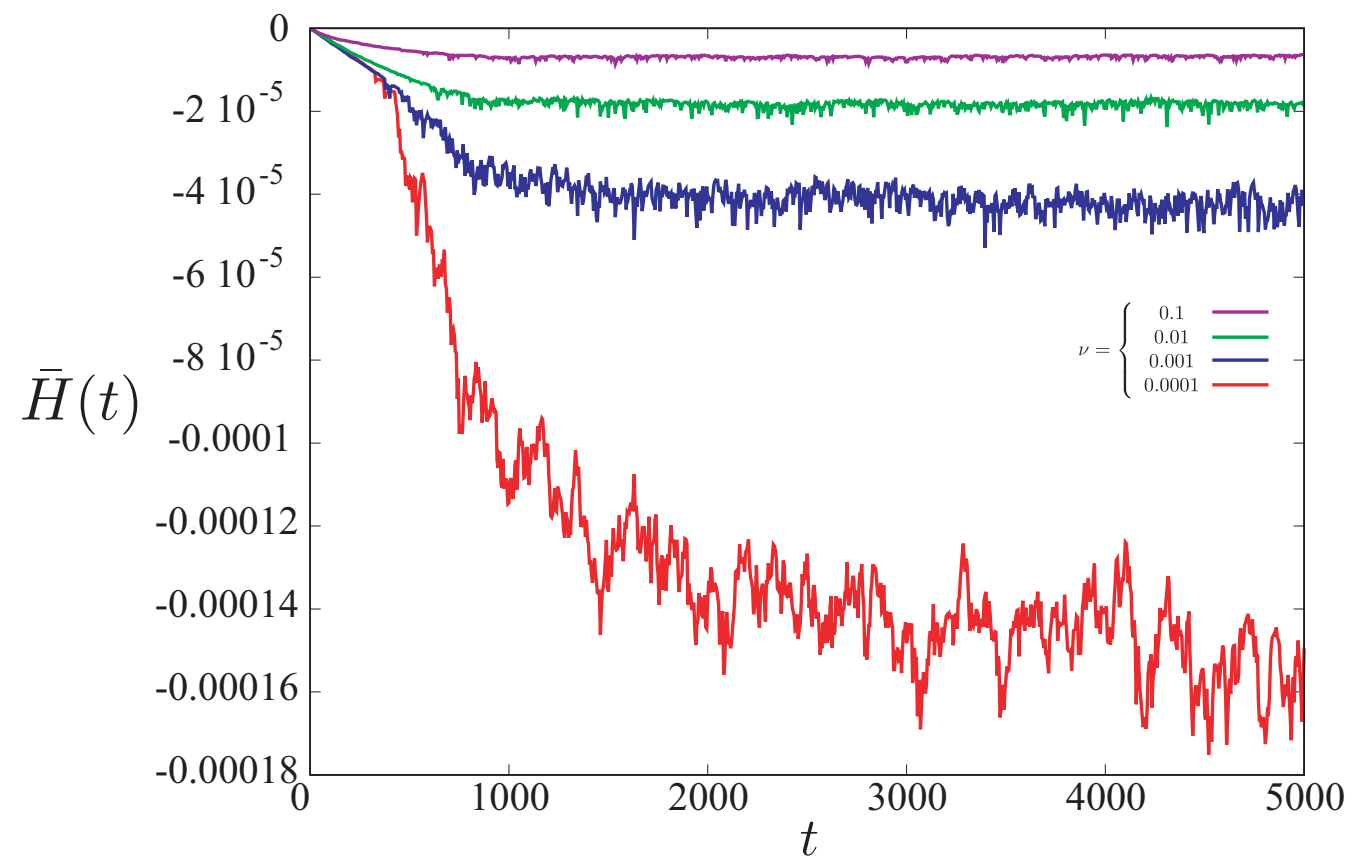
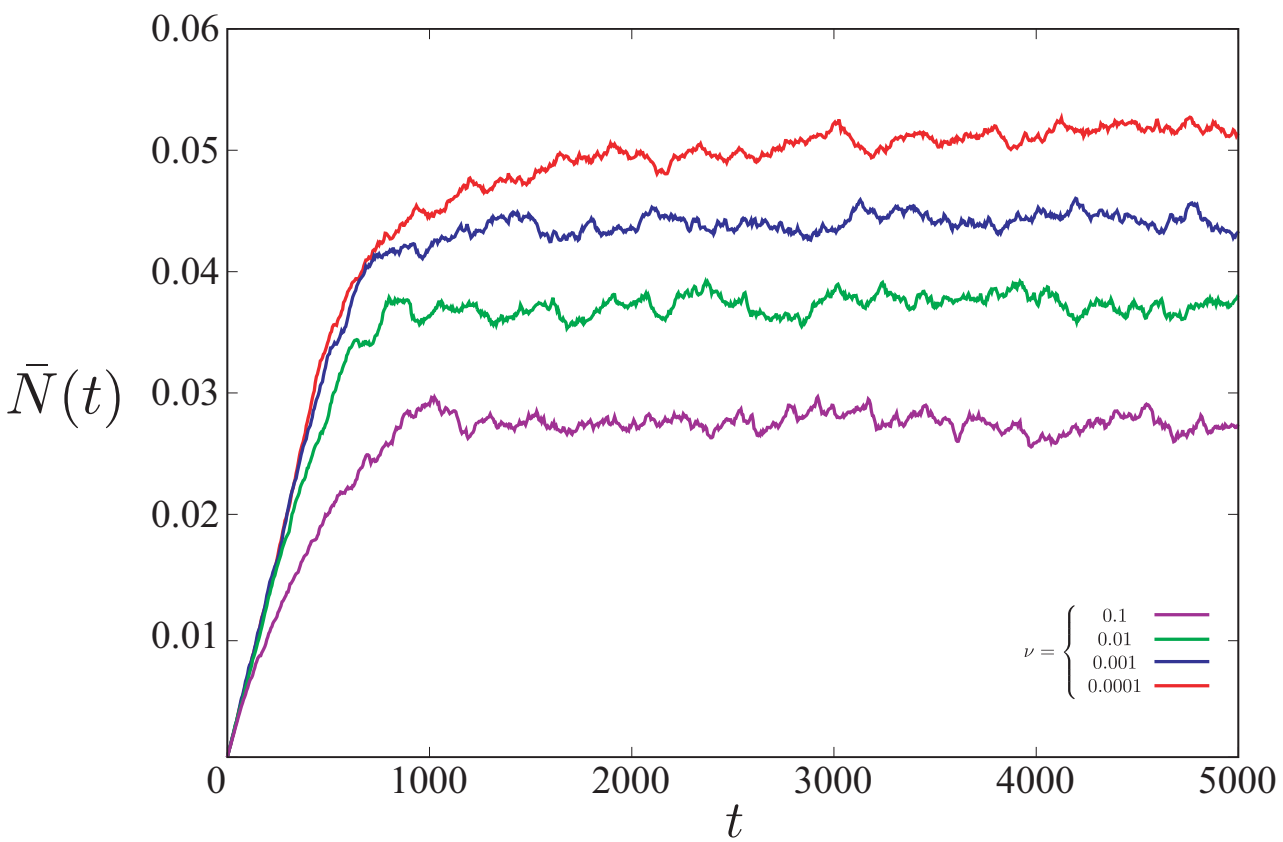
In higher dimensions?

- Similar investigation in 2D have been done

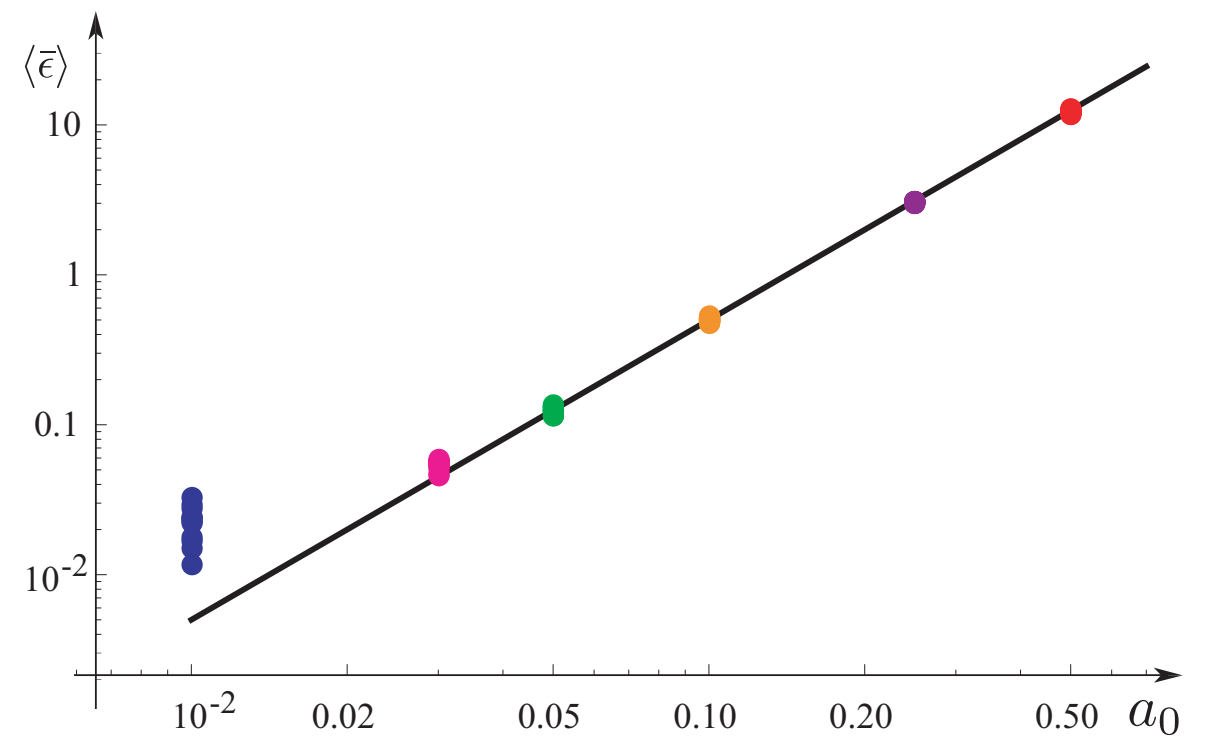
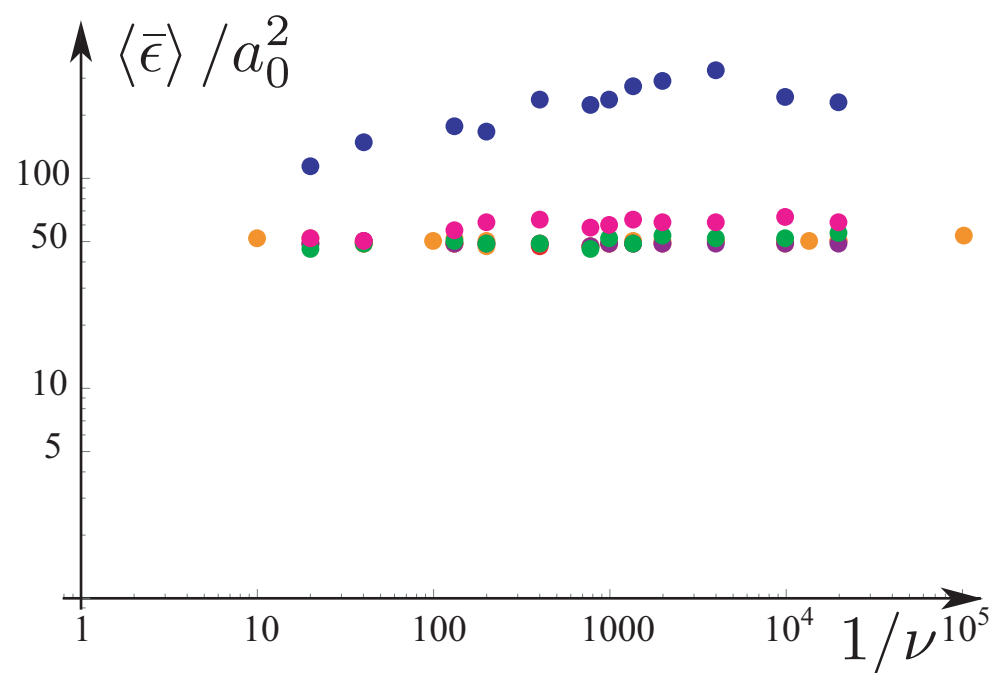
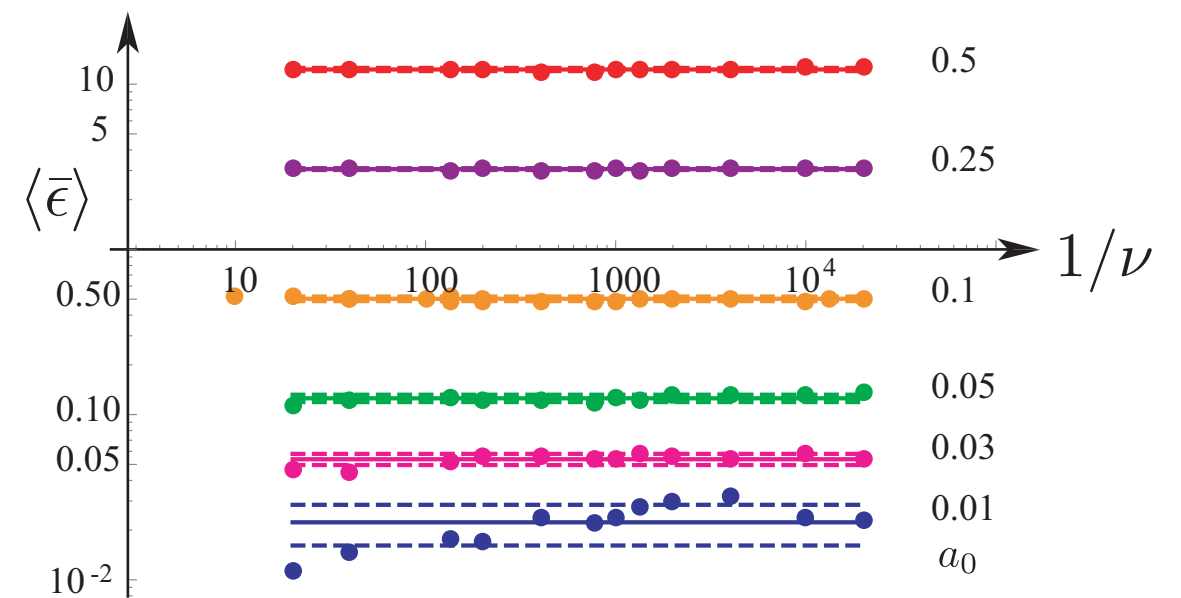
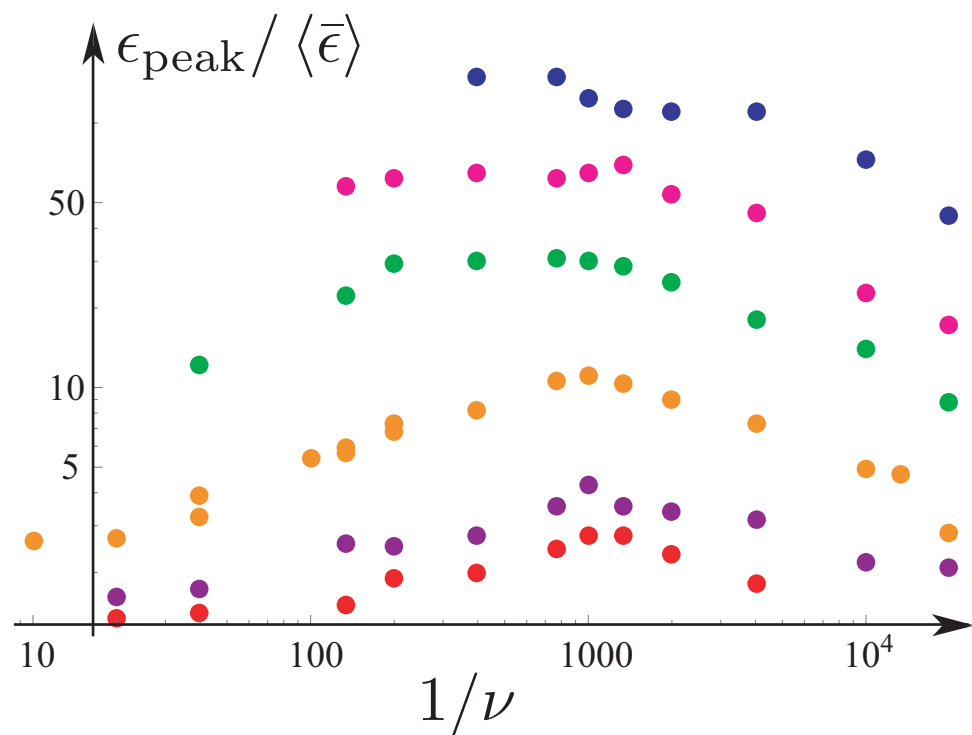


Turbulent state

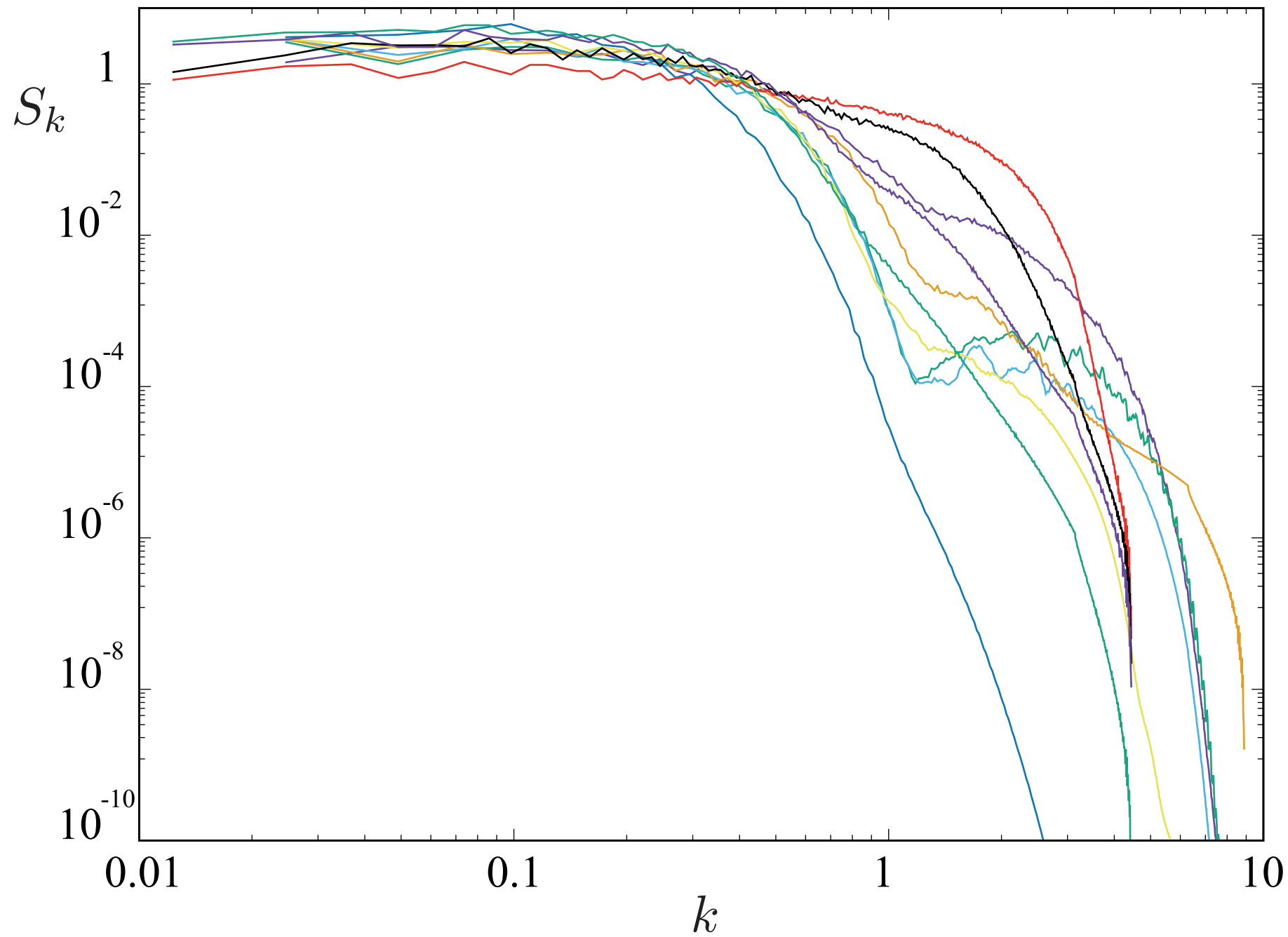




Influence of the forcing?



Flat spectrum?



$$S_k = A \frac{\alpha}{\epsilon^{1/3} g^{2/3}}$$

Conclusion

- « singularity » mediated turbulence (in the spirit of « defect » mediated (Coullet, Gil & Lega 1989) is observed (singularity cured by viscosity) in a version of the focusing NLS
- simple model where singularity in the inviscid limit is known. Mass « cascade »
- dissipation of mass concentrated in the collapses
- Kolmogorov like spectra are observed (needs additional condition for the exponent, different that those of the collapse and of the WTT, possibility of Phillips spectrum)
- singularity manifests through intermittency at high order